Confinement of Unimodal Distributions in High Dimension and an FKG-Gaussian Correlation Inequality

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Motivation: Unimodal Probability Distributions

Probability distributions take many forms. Which are simplest?

In 1-dimension, the unimodal distributions form a very nice class.



- Log-concave if $\log \rho(x)$ is concave.
- *M*-log-concave if for positive-definite *M* and all $x \in \mathbb{R}^N$:

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• Covariance bound [Brascamp-Lieb]:

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- μ inherits functional inequalities from γ_M [Bakry-Emery].
 - Spectral gap, isoperimetry, concentration.

Unimodality without Concentration

Unfortunately, some unimodal distributions do not behave so nicely. Consider the two-component Gaussian mixture

$$\frac{1}{2}\mathcal{N}(0,I_{\mathsf{N}})+\frac{1}{2}\mathcal{N}(0,4I_{\mathsf{N}}).$$

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This talk will provide one such tool.

• Results show confinement, which is weaker than concentration.

- Ginzburg–Landau Surfaces and Main Results
- Confinement from the Gaussian Correlation Inequality
- The FKG-Gaussian Correlation Inequality
- Putting it all together

Warmup: Gaussian free field (GFF) on locally finite graph G = (V, E): • Fix finite $\Lambda \subseteq V$, e.g. $[-L, \ldots, L]^d \subseteq \mathbb{Z}^d$. Set $\varphi(v) = 0$ for all $v \notin \Lambda$. Warmup: Gaussian free field (GFF) on locally finite graph G = (V, E):

- Fix finite $\Lambda \subseteq V$, e.g. $[-L, \ldots, L]^d \subseteq \mathbb{Z}^d$. Set $\varphi(v) = 0$ for all $v \notin \Lambda$.
- GFF is the random function $\phi: V \to \mathbb{R}$ with density:

$$d\mu_{G,\Lambda,GFF}(\varphi) = \frac{1}{Z_{G,\Lambda,GFF}} \exp\left(-\sum_{e=\{v,v'\}\in E} \frac{1}{2} |\varphi(v) - \varphi(v')|^2\right) \prod_{v\in\Lambda} d\varphi(v)$$

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- Models fluctuations of random interfaces.
- Lots of probabilistic interest, notably on \mathbb{Z}^2 (extreme values, LQG).

Discrete Gaussian Free Fields



Picture of GFF by Sam Watson. Here $\Lambda = [-L, \dots, L]^2 \subseteq \mathbb{Z}^2$.

Discrete Gaussian Free Fields and Electrical Networks

Well-known link between GFF and electrical networks:

• Let $R_{eff}(\cdot)$ be effective resistance on G. Then

$$\mathbb{E}^{\mu_{G,\Lambda,GFF}}[\phi(v)^2] = R_{eff}(v \leftrightarrow \partial \Lambda).$$

- More generally, $\mathbb{E}[(\phi(v) \phi(w))^2] = R_{eff}(v \leftrightarrow w).$
- $R_{eff}(v \leftrightarrow \infty) < \infty$ iff simple random walk on G is transient.

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Definition of effective resistance: the **energy** of $f : E(G) \rightarrow \mathbb{R}$ is

$$\mathcal{E}(f) = \sum_{e \in E} f(e)^2.$$

 $R_{eff}(v\leftrightarrow w)$ is the minimum energy of any v
ightarrow w unit flow.

Ginzburg-Landau Random Surfaces

More general Ginzburg-Landau surfaces use non-quadratic interactions:

$$\mathsf{d}\mu_{G,\Lambda,U}(\varphi) \equiv \frac{1}{Z_{G,\Lambda,U}} \exp\left(-\sum_{e \in E} \frac{U}{\nabla \varphi(e)}\right) \prod_{v \in \Lambda} \mathsf{d}\varphi(v).$$

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- First rigorous study in [Brascamp-Lieb-Lebowitz 1975].
- Names: "Ginzburg–Landau", " $\nabla \phi$ ", "anharmonic crystal".
- Dynamics, large deviations, fluctuations, Z-valued analogs,... [Funaki-Spohn 97, Naddaf-Spencer 97, Deuschel-Giacomin-Ioffe 00, Sheffield 03, Miller 11, Armstrong-Dario 22, Armstrong-Wu 23...].

Localization

We will consider the question of localization.

Question

Are fluctuations of $\varphi(v_0)$ stochastically bounded on large domains $\Lambda \uparrow V$? If so, we say the model is **localized**. Otherwise **delocalized**.

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 - One can even take this as the definition of localization.
- GFF on \mathbb{Z}^d localizes iff $d \geq 3$.
 - Equivalent to transience/recurrence since

$$\mathbb{E}^{\mu_{G,\Lambda,GFF}}[\phi(v)^2] = R_{eff}(v \leftrightarrow \partial \Lambda).$$

• On $[-L, \ldots, L]^2 \subseteq \mathbb{Z}^2$, one has $\mathbb{E}[\phi(\vec{0})^2] \approx \log L$.

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- On $[-L,\ldots,L]^2\subseteq \mathbb{Z}^2,$ one has $\mathbb{E}[\phi(\vec{0})^2]\approx \log L.$
- Conjecture of [Brascamp-Lieb-Lebowitz 1975]: localization is determined by the **geometry** of *G*, not the potential *U*.
 - Proved delocalization for very general $U \in C^2(\mathbb{R})$ on \mathbb{Z}^2 .

Localization of Ginzburg-Landau Random Surfaces

Localization is known for various U (often focused on lattices):

- Strongly convex potentials with $\inf_{x \in \mathbb{R}} U''(x) \ge c > 0.$
 - $\mu_{G,\Lambda,U}$ is dominated by GFF [Brascamp-Lieb-Lebowitz 75].

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- U(x) = |x| using infrared bounds [Bricmont-Fontaine-Lebowitz 82].
- Mildly non-convex U via renormalization [Cotar-Deuschel-Muller 09 & 12, ABKM 16 & 19, Hilger 16 & 20].

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- e^{-U(x)} is a mixture of centered Gaussians (will explain soon)
 [Biskup-Kotecky 07, Biskup-Spohn 11, Brydges-Spencer 12, Ye 19,...].
- [Magazinov-Peled 22]: convex U with U''(x) > 0 for a.e. x.

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- [Magazinov-Peled 22]: convex U with U''(x) > 0 for a.e. x.
- Still open for Hammock potential U(x) = ∞ · 1_{|x|>1}. This gives a uniformly random 1-Lipschitz φ : V → ℝ.

11/44

We prove localization for monotone potentials.

Definition ((α, ε)-monotonicity)

U is (α, ε) -monotone if it is even, increasing on \mathbb{R}^+ , and $U'(x) \ge \min\left(\varepsilon x, \frac{1+\alpha}{x}\right)$ for all points of differentiability $x \ge 0$.

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Theorem (Localization for (α, ε) -monotone U)

Let G be transient, and U be (α, ϵ) -monotone for $\alpha > 2$. Then $\mathbb{P}^{\mu_{G,\Lambda,U}}[|\phi(v_0)| \ge t] \le O(t^{-\alpha})$ uniformly in $\Lambda \subseteq V$, for any $v_0 \in V$.

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- Proof will be based on unimodality of $\mu_{G,\Lambda,U}$.
- U_e can depend on edge e, as long as (α, ε) are uniform.
- If G is transient and transitive, $\varphi(v_0)$ is tight even for $\alpha = \epsilon$.
 - \approx minimal condition for $\int_{\mathbb{R}} e^{-U(x)} dx < \infty$ so $Z_{G,\Lambda,U} < \infty$.

Extreme Values of the Field

These bounds are often sharp enough to understand $\max_{v \in \Lambda} |\phi(v)|$.

Theorem (Extreme Values from Polynomial Bounds)

Let U be (α, ε) -monotone with $\sup_{x \ge 1} |U(x) - (1 + \alpha) \log x| < \infty$ and $\alpha > 2$. As $\Lambda \subseteq \mathbb{Z}^d$ varies for $d \ge 3$, the laws of

$$|\Lambda|^{-\frac{1}{2d\alpha}} \max_{v \in \Lambda} |\varphi(v)|$$

are tight in $(0, \infty)$, i.e. stochastically bounded away from 0 and ∞ .

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- Upper bound: Markov with extra tricks to get 2dα in the exponent. (Split Z^d into 2d transient subgraphs containing the vertex v...)
- Similar condition for stretched exponential tails.
 - Monotonicity condition: $U'(x) \ge \min(\epsilon x, \epsilon x^{\beta-1})$, for $\beta \in (0, 2]$.

General Statement without Graphs

The graph structure is irrelevant in the main result!

Let U be (α, ε) -monotone for $\alpha > 2$, and $\ell_1, \ldots, \ell_j : \mathbb{R}^d \to \mathbb{R}$ be linear.

Choose $\varphi, \widetilde{\varphi} \in \mathbb{R}^d$ from densities:

$$\begin{split} \phi &\sim \exp\Big(\sum_{i=1}^{j} - \mathcal{U}(\ell_{i}(\phi))\Big) d\phi/Z_{\vec{\ell},\mathcal{U}}, \\ \widetilde{\phi} &\sim \exp\Big(\sum_{i=1}^{j} - \ell_{i}(\widetilde{\phi})^{2}\Big) d\widetilde{\phi}/Z_{\vec{\ell},\textit{GFF}}. \end{split}$$

Fix any other linear function $\ell_* : \mathbb{R}^d \to \mathbb{R}$. Then $\ell_*(\varphi)$ is bounded on the same scale as the centered Gaussian $\ell_*(\widetilde{\varphi})$, with α -power tails.

• Recovering GFF/Ginzburg-Landau: set $\ell_e(\phi) = \phi(v) - \phi(v')$.

Preview of the Proof

The proof has two core components:

• Handle the case that U = V takes the form:

$$e^{-V(x)} = \int_0^\infty \frac{e^{-x^2/2\xi^2}}{\xi\sqrt{2\pi}} \mathrm{d}
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- The overall model will be a mixture of Gaussian processes, e.g. GFFs with edge weights [Biskup-Kotecky 07].
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- [Biskup-Spohn 11]: can have phase coexistence pprox non-concentration.
- Reduce to this case using the FKG-Gaussian correlation inequality. Amounts to domination by mixtures of centered Gaussians.
 - Dominating Gaussian mixtures must have special structure.
 - Perfectly suited for products of 1-dimensional functions.

- Results on Ginzburg–Landau Model
- Confinement from the Gaussian Correlation Inequality
- The FKG-Gaussian Correlation Inequality
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Royen's Gaussian Correlation Inequality

Theorem (Royen 2014)

Let γ be a centered Gaussian measure on \mathbb{R}^d , and $K_1, K_2 \subseteq \mathbb{R}^d$ symmetric convex sets (i.e. $K_i = -K_i$). Then 1_{K_1} and 1_{K_2} have non-negative correlation under γ , i.e.

 $\gamma(K_1 \cap K_2) \geq \gamma(K_1)\gamma(K_2).$



Theorem (Royen 2014)

For γ centered Gaussian on \mathbb{R}^d , and $K_1, K_2 \subseteq \mathbb{R}^d$ symmetric convex sets:

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History (see 2017 Quanta article):

- Conjectured by [Dunnet-Sobel 55], [Gupta-Eaton-Perlman-Savage-Sobel 72].
- [Khatri 67, Sidak 67, Pitts 77, Schechtman-Schlumprecht-Zinn 98, Hargé 99]: special cases such as \mathbb{R}^2 .
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Proof idea: for x, $y \stackrel{i.i.d.}{\sim} \gamma$, equivalent to

$$\mathbb{P}[x \in K_1 \land x \in K_2] \geq \mathbb{P}[x \in K_1, y \in K_2].$$

Royen showed $f(t) = \mathbb{P}[x \in K_1 \land \sqrt{1-t}x + \sqrt{t}y \in K_2]$ is decreasing.

18/44

Symmetric Quasi-Concave Functions

Definition

- $f: \mathbb{R}^N \to \mathbb{R}$ is symmetric quasi-concave (SQC) if:
 - f(x) = f(-x) for all $x \in \mathbb{R}^N$.
 - All super-level sets $\{x \in \mathbb{R}^N : f(x) \ge \lambda\}$ are convex.



GCI: if $K_1, K_2 \subseteq \mathbb{R}^d$ are symmetric convex, then

 $\gamma(\mathit{K}_1\cap \mathit{K}_2)\geq \gamma(\mathit{K}_1)\gamma(\mathit{K}_2).$

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If $K_1, \ldots, K_{m+1} \subseteq \mathbb{R}^d$ are symmetric convex:

 $\gamma(\mathcal{K}_1 \cap \cdots \cap \mathcal{K}_{m+1}) \geq \gamma(\mathcal{K}_1 \cap \cdots \cap \mathcal{K}_m) \cdot \gamma(\mathcal{K}_{m+1}),$

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By level sets, if $f_1, \ldots, f_{m+1} : \mathbb{R}^d \to \mathbb{R}^+$ are symmetric quasi-concave, $\mathbb{E}^{\gamma}[f_1 f_2 \ldots f_{m+1}] \ge \mathbb{E}^{\gamma}[f_1 f_2 \ldots f_m] \cdot \mathbb{E}^{\gamma}[f_{m+1}].$

(Products of SQC functions need not be SQC, hence the middle step.)

If $f_1, \ldots, f_{m+1} : \mathbb{R}^d \to \mathbb{R}^+$ are symmetric quasi-concave,

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Suppose γ is centered Gaussian and $\frac{dv}{d\gamma} = f_1 f_2 \dots f_m$ is a product of SQC functions. Then

$$\nu(\mathcal{K}) = \mathbb{E}^{\gamma} \left[\frac{d\nu}{d\gamma} \cdot \mathbf{1}_{\mathcal{K}} \right] \stackrel{\mathsf{GCI}}{\geq} \mathbb{E}^{\gamma} \left[\frac{d\nu}{d\gamma} \right] \cdot \gamma(\mathcal{K}) = \gamma(\mathcal{K})$$

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for symmetric convex K.

This is a form of **Gaussian domination**. We say $\nu \leq_{con} \gamma$.

Definition

 $\nu \preceq_{con} \gamma$ if $\gamma(K) \le \nu(K)$ for all symmetric convex sets K.

Consequence: localization on all transient G if $U'(x) \ge \varepsilon x$ for all $x \ge 0$.

Why? Domination by rescaled GFF γ_{ε} with potential $U_{\varepsilon}(x) = \varepsilon x^2/2$.

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Why? Domination by rescaled GFF γ_{ε} with potential $U_{\varepsilon}(x) = \varepsilon x^2/2$.

The Radon-Nikodym derivative is a product:

$$\frac{\mathrm{d}\mu_{G,\Lambda,U}}{\mathrm{d}\gamma_{\mathcal{E}}} \propto \prod_{e \in E(G)} W(\nabla \varphi(e)), \quad \text{with} \quad W(x) = e^{-U(x) + U_{\mathcal{E}}(x)}.$$

Consequence: localization on all transient G if $U'(x) \ge \varepsilon x$ for all $x \ge 0$.

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Conclusion: $\mu_{G,\Lambda,U} \preceq_{con} \gamma_{\varepsilon}$. Localization on all transient G.

This method requires that $U(x) \ge \Omega(x^2)$. μ must have subgaussian tails to be dominated by a single Gaussian.

- Results on Ginzburg–Landau Model
- Confinement from the Gaussian Correlation Inequality

• Putting it all together

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Does GCI extend to **mixtures** of centered Gaussians? If so, we could use them as the dominating measures and have more flexibility.

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If the Gaussians have comparable covariance, GCI extends!!

Theorem

Let $\Sigma_1 \succeq_{PSD} \Sigma_2 \succeq_{PSD} 0$ be symmetric matrices. Let $d\gamma_1(x) \propto e^{-\langle x, \Sigma_1 x \rangle}$ and $d\gamma_2(x) \propto e^{-\langle x, \Sigma_2 x \rangle}$. Then GCI holds for $\mu = p\gamma_1 + (1-p)\gamma_2$:

 $\mu({\cal K}\cap{\cal K}')\geq \mu({\cal K})\mu({\cal K}')$

for any symmetric convex sets K, K' and $0 \le p \le 1$.



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- Can show $\gamma_1(K) \ge \gamma_2(K)$ (by GCI or otherwise). Thus $q \ge p$:

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$$\widetilde{\mu}(\mathcal{K}') = q\widetilde{\gamma}_{1}(\mathcal{K}') + (1-q)\widetilde{\gamma}_{2}(\mathcal{K}') \stackrel{GCI}{\geq} q\gamma_{1}(\mathcal{K}') + (1-q)\gamma_{2}(\mathcal{K}')$$

$$\stackrel{Rearr}{\geq} p\gamma_{1}(\mathcal{K}') + (1-p)\gamma_{2}(\mathcal{K}') = \mu(\mathcal{K}').$$

GCI For Totally Ordered Gaussian Mixtures

We can generalize further! Suppose:

- $\mu = p_1 \gamma_1 + \dots + p_j \gamma_j$, with totally ordered inverse covariances $\Sigma_1 \succeq_{PSD} \Sigma_2 \succeq_{PSD} \dots \succeq_{PSD} \Sigma_j$.
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An analogous proof shows $\widetilde{\mu}(K') \ge \mu(K')$ for any symmetric convex K':

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• We similarly get:

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We need these two functions on $\{1, 2, ..., j\}$ to be positively correlated with respect to the probability measure $\mathbb{P}[i] = p_i$.

This is the rearrangement inequality, a special case of FKG.

Log-Supermodular Gaussian Mixtures (LSGM)

Example $(2 \times 2 \text{ Lattice})$

Let
$$p_{1,1}p_{2,2} \ge p_{1,2}p_{2,1}$$
. Suppose $d\gamma_{i,j}(x) \propto e^{-\langle x, \Sigma_{i,j} x \rangle}$ with:

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Definition (Log-Supermodular Gaussian Mixture)

An LSGM on \mathbb{R}^n is a Gaussian mixture $\Gamma_{\nu,\Sigma} = \int \gamma_{\xi} d\nu(\xi)$. such that: • $d\nu(x) = f(x)dx$ is log-supermodular on \mathbb{R}^k_+ :

$$f(\xi)f(\xi')\leq f(\xi\wedge\xi')f(\xiee\xi'),\quad orall\,\xi,\xi'\in\mathbb{R}^k_+.$$

•
$$d\gamma_{\xi}(x) \propto e^{-\langle x, \Sigma(\xi)x \rangle}$$
, for some $\Sigma : \mathbb{R}^{k}_{+} \to \mathcal{S}^{n}_{+}$.

• Σ is order-reversing from \leq_{coord} to \leq_{PSD} .

Theorem (FKG-GCI)

For any LSGM $\Gamma_{\nu,\Sigma}$ and symmetric convex K_1, K_2 : $\Gamma_{\nu,\Sigma}(K_1 \cap K_2) \ge \Gamma_{\nu,\Sigma}(K_1)\Gamma_{\nu,\Sigma}(K_2).$

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Corollary (Domination by LSGM Gaussian Mixtures)

For any LSGM $\Gamma_{v,\Sigma}$, suppose for symmetric quasi-concave f_1, \ldots, f_m :

$$\mathrm{d}\widetilde{\Gamma}(x) = f_1(x)f_2(x)\ldots f_m(x) \,\mathrm{d}\Gamma_{\mathbf{v},\boldsymbol{\Sigma}}(x).$$

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The FKG-Gaussian Correlation Inequality

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For Ginzburg-Landau, it remains to express $\mu_{G,\Lambda,U}$ as Γ above.

Although GCI is quite elegant, the proof is miraculous and took decades to find. For the purposes of this talk, it can be bypassed.

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"Slab GCI" (Sidak 68): let K, K' be symmetric convex sets with

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a symmetric slab. Then for centered Gaussian γ ,

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Instead of \leq_{con} , one may consider the weaker relation

if $v(K) \ge \mu(K)$ for all symmetric slabs K. Then slab GCI yields a slab-FKG-GCI inequality. This will recover all our localization results, since they concern 1-dimensional projections of φ .

M. Sellke

- Ginzburg–Landau Surfaces and Main Results
- Confinement from the Gaussian Correlation Inequality
- The FKG-Gaussian Correlation Inequality
- Putting it all together

Dominating LSGMs will be $\mu_{G,\Lambda,V}$ where V takes the form:

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$$d\mu_{G,\vec{\xi},GFF}(\varphi) = \frac{1}{Z_{G,GFF}} \exp\left(-\sum_{e \in E} \frac{1}{2\xi_e^2} \cdot |\nabla\varphi(e)|^2\right) \prod_{v \in V} d\varphi(v).$$

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Encoding: inverse covariance $\Sigma(\vec{\xi})$ given by

$$\left\langle \varphi, \Sigma(\vec{\xi})\varphi \right\rangle = \sum_{e \in E(G)} (\varphi(v) - \varphi(v'))^2 / \xi_e^2.$$

• Clearly Σ is order-reversing from \leq_{coord} to \leq_{PSD} .

Potential $e^{-V(x)} = \int_0^\infty \frac{e^{-x^2/2\xi^2}}{\xi\sqrt{2\pi}} d\rho(\xi)$. Encoding: inverse covariance

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The Ginzburg-Landau measure is a mixture of these GFFs:

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Mixing measure $d\nu(\vec{\xi})$ is **not** a product. It gains a factor $det(\Sigma(\xi))^{-1/2}$.

• Elementary fact: if A, B, $C \succeq_{PSD} 0$, then

 $\det(A)\det(A+B+C) \leq \det(A+B)\det(A+C).$

• This yields log-supermodularity, since $\Sigma(\vec{\xi}) = \sum_{e \in E} F(\xi_e)$ is additive.

If $U'(x) \ge V'(x)$ on \mathbb{R}_+ , the Radon–Nikodym derivative

$$\frac{\mathrm{d}\mu_{G,\Lambda,U}}{\mathrm{d}\mu_{G,\Lambda,V}} \propto \prod_{e \in E} e^{-U(\nabla \varphi(e)) + V(\nabla \varphi(e))}$$

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Further, $\mu_{G,\Lambda,V}$ is dominated by the "naive product" LSGM:

$$\mu_{G,\Lambda,V} \preceq_{\mathsf{con}} \int \mu_{G,\vec{\xi},GFF} \prod_{e \in E} \mathsf{d}\rho(\xi_e).$$

Indeed the presence of $G \setminus \{e\}$ only makes ξ_e smaller, so one gets stochastic domination of $\vec{\xi}$ by $\prod_{e \in E} \rho(\xi_e)$. This reduces localization to studying GFFs with IID edge resistances.

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Lemma

There exist potentials $V(\rho)$ in centered Gaussian mixture form such that:

$$\begin{array}{l} \bullet \quad V'(x) \leq \min\left(\epsilon x, \frac{1+\epsilon}{x}\right), \quad \forall x \geq 0. \\ \bullet \quad V'(x) \leq \min\left(\epsilon x, \frac{1+\alpha}{x}\right) \text{ and } \rho([t,\infty)) \leq O(t^{-\alpha}), \quad \forall t \geq 0. \\ \bullet \quad V'(x) \leq \min\left(\epsilon x, \epsilon x^{\beta-1}\right) \text{ and } \rho([t,\infty)) \leq e^{-\Omega(t^{\beta})}, \quad \forall t \geq 0. \end{array}$$

In each case, $U' \ge V'$ if U is correspondingly monotone.

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In each case, $U' \ge V'$ if U is correspondingly monotone.

Proof Idea: Explicit construction. Match tail of ρ to the decay rate.

$$\ \, \mathfrak{o}([t,\infty)) \asymp e^{-t^{\beta}}.$$

Fix $\alpha > 2$ and transient G. Let U be (α, ε) -monotone. Then

$$\mathbb{P}^{\mu_{G,\Lambda,U}}[|\varphi(v)| \geq t] \leq O(t^{-lpha})$$

holds uniformly in Λ for any transient G.

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Proof: The variance of $\varphi(v)$ after $\vec{\xi}$ -weighting is a weighted effective resistance $R_{eff}^{(\xi)}(v \leftrightarrow \infty)$. So we must bound the tail of $R_{eff}^{(\xi)}(v \leftrightarrow \infty)$.

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$$\mathbb{P}^{\mu_{G,\Lambda,U}}[|\phi(v)| \geq t] \leq O(t^{-lpha})$$

holds uniformly in Λ for any transient G.

We'll compare to GFF with IID $\xi_e \sim \rho$, with $\rho([t, \infty)) \leq O(t^{-\alpha})$.

Proof: The variance of $\varphi(v)$ after $\vec{\xi}$ -weighting is a weighted effective resistance $R_{eff}^{(\xi)}(v \leftrightarrow \infty)$. So we must bound the tail of $R_{eff}^{(\xi)}(v \leftrightarrow \infty)$.

Consider the energy-minimizing unit flow $v \to \infty$ in the **unweighted** *G*. Its weighted energy is random and upper bounds $R_{eff}^{(\xi)}(v \leftrightarrow \infty)$. This is $\sum_e a_e \xi_e^2$, where $\sum_e a_e = R_{eff}(v \leftrightarrow \infty) < \infty$. Now use e.g. Jensen. \Box

Suppose U is $(\varepsilon, \varepsilon)$ -monotone, and p-bond percolation on G has transient infinite cluster for $p \in [1 - \delta, 1]$. Then Law $(\varphi(v))$ is tight as $\Lambda \uparrow \infty$.

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By definition, these edges form a transient infinite cluster \mathcal{C} .

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Let $w \in C$ be the closest point to v. Then both $R_{eff}^{(\xi)}(v \leftrightarrow w)$ and $R_{eff}^{(\xi)}(w \leftrightarrow \partial \Lambda)$ are tight. Hence $R_{eff}^{(\xi)}(v \leftrightarrow \partial \Lambda)$ is also tight.

A Remark on Monotonicity

We used that if $U'(x) \ge V'(x)$ for all $x \ge 0$, then

 $\mu_{G,\Lambda,U} \preceq_{\operatorname{con}} \mu_{G,\Lambda,V}.$

This holds edge-by-edge. Hence if $\beta_e \ge 1$ for each $e \in E(G)$, then

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I.e. for Gaussian-mixture potentials V, "stiffening the springs" improves confinement. [BLL75] gives a **counterexample** for

$$V(x) = x^2 + \varepsilon x^4.$$

IV. STIFFENING THE SPRINGS DOES NOT NECESSARILY DECREASE $\langle x_6^2 \rangle$

Consider the following case with three particles, i.e.

 $\begin{array}{l} H=v(x)+v(x-y)+v(y-z)+v(z)+\alpha v(x-z) \mbox{ and } x_0=y. \mbox{ Let } v(x)=x^2+\varepsilon x^4,\\ \varepsilon>0. \mbox{ We want to show that in creasing α from 0 can decrease $<x_0^2>. \mbox{ Let } g_{\infty}$ (resp. g_0) be $<x_0^2>$ for $\alpha=\infty$ (resp. $\alpha=0$). \mbox{ Then } g_{\infty}=2$ $\int y^2 G(y)/[G(y)$ and $g_0=\int y^2 F(y)/[F(y)$, with $G(y)=\exp(-2v(y)$)$ and $F(y)=R(y)^2$ where $R=\exp(-v)*\exp(-v)$. A simple calculation shows that for the pure harmonic case $(\varepsilon=0)$, $g_0=g_{\infty}=1/2$. When $\varepsilon>0$ it is impossible to calculate the integrals, but it is possible to calculate $g_i=dg_i/de_{i=0}$. One finds that $g_{\infty}=-3/4$ and $g_0=-9/8$. Thus, for small, positive ξ, $g_{\infty}>g_0$, which is the contradiction we wished to demonstrate. \end{tabular}$

- Ginzburg–Landau Surfaces and Main Results
- Confinement from the Gaussian Correlation Inequality
- The FKG-Gaussian Correlation Inequality
- Putting it all together

• Another Application: the Fröhlich Polaron

Let $d\mathbb{Q}(B)$ be the law of 3-dimensional Brownian motion.

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Given coupling strength $\alpha \gg 1$ and time-horizon $T \gg \alpha$, the **Polaron** path measure $\widehat{\mathbb{Q}}_{\alpha,T}$ is the reweighted law on paths $B : [0, T] \rightarrow \mathbb{R}^3$:

$$\mathrm{d}\widehat{\mathbb{Q}}_{\alpha,T}(\mathsf{B}) \equiv \frac{1}{Z_{\alpha,T}} \exp\left(\alpha \int_{0}^{T} \int_{0}^{T} \frac{e^{-|t-s|}}{\|\mathsf{B}_{t}-\mathsf{B}_{s}\|} \, \mathrm{d}t \, \mathrm{d}s\right) \mathrm{d}\mathbb{Q}(\mathsf{B}).$$

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Obtained by Feynman's path integral applied to a quantum operator (modeling an electron in crystal). The "effective mass" is

$$m_{eff}(\alpha) = \lim_{T \to \infty} \frac{3T}{\mathbb{E}^{\widehat{\mathbb{Q}}_{\alpha,T}} \|\mathsf{B}_{T}\|^{2}} \stackrel{?}{\approx} C_{*} \alpha^{4}.$$

[Fröhlich 37, Landau-Pekar 48, Feynman 55, Lieb 77, Donsker-Varadhan 83, Spohn 87, Lieb-Thomas 97, Lieb-Seiringer 17, Mukherjee-Varadhan 18 & 20, Dybalski-Spohn 20, Betz-Polzer 21 & 22, Brooks-Seiringer 22, S 22]

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In fact, the Coulomb interaction is a mixture of centered Gaussians:

$$\frac{1}{x} = \sqrt{2/\pi} \int_0^\infty e^{-u^2 x^2/2} \mathrm{d}u.$$

Hence $\widehat{\mathbb{Q}}_{\alpha, \mathcal{T}}$ is an (infinite dimensional) LSGM and obeys GCI.

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Now the mixture comes **inside** the exponent. The resulting Gaussian mixture representation of $\widehat{\mathbb{Q}}_{\alpha,T}$ is indexed by a deformed Poisson process on weighted time-intervals ([*s*, *t*]; *u*).

Theorem (Mukherjee-Varadhan 20)

The Polaron path measure $\widehat{\mathbb{Q}}_{\alpha,\,T}$ has a mixture-of-Gaussian representation

$$\widehat{\mathbb{Q}}_{\alpha,T}(B_{[0,T]}) = \int \mathsf{Q}_{\xi}(B_{[0,T]}) \ \widehat{\Theta}_{\alpha,T}(\mathsf{d}\xi).$$

Here $\xi = \{([s_i, t_i], u_i)\}_{i=1}^n$ is a point process of weighted intervals, and $dQ_{\xi}(B_{[0,T]}) \propto e^{-\sum_{i=1}^n u_i^2 ||B(t_i) - B(s_i)||^2} d\mathbb{Q}(B_{[0,T]}).$
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- Ergodic limit $(\widehat{\Theta}_{\alpha, \mathcal{T}}, \widehat{\mathbb{Q}}_{\alpha, \mathcal{T}}) \to (\widehat{\Theta}_{\alpha, \infty}, \widehat{\mathbb{Q}}_{\alpha, \infty})$ as $\mathcal{T} \to \infty$.
- Functional CLT for Q_{α,∞} [Mukherjee-Varadhan 20, Betz-Polzer 21].
 Rigorizes path integral connection [Spohn 87, Dybalski-Spohn 20].
- Applied to show $m_{eff}(lpha)\gtrsim lpha^{2/5}$ [Betz-Polzer 22].
- [S 22,Brooks-Seiringer 22]: $\frac{\alpha^4}{C \log(\alpha)^6} \le m_{eff}(\alpha) \le C_* \alpha^4 + \alpha^{4-\epsilon}.$

The Fröhlich Polaron

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Theorem (Bazaes-Mukherjee-S-Varadhan 24; predicted in Landau-Pekar 1948)

 $\textit{If $T^{0.1} \geq \alpha \gg 1$, then $\mathbb{E}^{\widehat{\mathbb{Q}}_{\alpha, T}} \| B_T \|^2 \leq O(T\alpha^{-4})$. i.e. $m_{eff}(\alpha) \geq \Omega(\alpha^4)$.}$

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Here FKG-GCI is "technical glue" for monotonicity and local-to-global arguments. Some direct consequences:

•
$$\alpha \mapsto m_{eff}(\alpha)$$
 is strictly increasing.

- $\mathbb{E}^{\widehat{\mathbb{Q}}_{\alpha,\tau_1+\tau_2}} \|B_{\mathcal{T}_1+\mathcal{T}_2}\|^2 \leq \mathbb{E}^{\widehat{\mathbb{Q}}_{\alpha,\tau_1}} \|B_{\mathcal{T}_1}\|^2 + \mathbb{E}^{\widehat{\mathbb{Q}}_{\alpha,\tau_2}} \|B_{\mathcal{T}_2}\|^2.$
- Universality: can replace 1/x by "more attractive" potentials.

Idea: find many "interval chains" $[s_1, t_1], [s_2, t_2], \ldots$ with $t_k \approx s_{k+1}$.

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- Sharp bounds for Ginzburg-Landau and Polaron models:

$$d\mu_{G,\Lambda,U}(\varphi) \equiv \frac{1}{Z_{G,\Lambda,U}} \exp\left(-\sum_{e \in E(G)} U(|\nabla \varphi(e)|)\right) \prod_{v \in \Lambda} d\varphi(v),$$
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• Anything similar without origin-symmetry? E.g. non-zero tilts.